# Numerical Study of an Indicator Function for Front-Tracking Methods 

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Received 9 May 2022; Revised 17 June 2022; Accepted 4 July 2022; Published 31 July 2022
Academic Editor: Zhengbiao Peng
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#### Abstract

In this paper, we present a detailed derivation and numerical investigation of an indicator function for front-tracking methods. We use the discrete Dirac delta function to construct an indicator function from a set of Lagrangian points and solve the resulting discrete Poisson equation with the zero Dirichlet boundary condition using an iterative method. We present several computational tests to investigate the effect of parameters such as distance between points, uniformity of the distance, and types of the Dirac delta functions on the indicator function.


## 1. Introduction

Dirac delta function plays an important role in various scientific and industrial problems. One of them is on the numerical simulations of fluid flows including multiphase flows and fluid solid interaction (FSI) problems. The immersed boundary methods (IBM) are on the front burner for the numerical simulations of these problems and one of the key points lay on the methods to accurately separate the moving multiphase or solid boundaries. The phase field method (PHM) [1-4], level set method (LSM) [5-7], and volume of fluid methods (VOF) [8-10], express the different phases using particular scalar functions. These functions are designed to be advected by fluid flows and one phase changes to another one continuously to avoid numerical stability problem. It should be noted that the interphase is diffused and implicitly captured in the three methods. However, the immersed boundary method is different for its explicit boundary expression. Unlike PHM, LSM, and VOF, distinguish of different phase can be denoted directly by scalar phase
functions; in IBM, a so-called indicator function needs to be constructed for these phase distinctions indirectly. In this process, if a Dirac delta function is used for the construction of the indicator function, the diffused interaction forms as PHM, LSM, and VOF. From this point of view, these four methods are all belonged to diffused models in simulations of multiphase flows. In IBM, the fluid interface is explicitly represented by several Lagrangian points and the fluid flow is solved on a stationary Eulerian mesh. Furthermore, when solving Richards' equation [11, 12] for layered soils, the layers of soils can be expressed using an indicator function.

Multiphase fluid flows are important in various scientific and industrial problems. There are many mathematical models and numerical methods for the multiphase fluid flows such as the immersed boundary method (IBM) [13-16], level set method [5-7], phase-field method [1-4], volume of the fluid method $[9,10]$, and lattice Boltzmann method [17-20], to name a few.

In the front-tracking method [21] such as IBM, the fluid-fluid interface is represented by a Lagrangian grid and
the flow is solved on a stationary Eulerian mesh. The Lagrangian points are represented by a set of marker points and move with the fluid flow defined at the Eulerian grid. To represent viscosity or density difference between different fluids, and indicator function is calculated based on the marker points [21-23]. Therefore, the indicator function has an important role in not only theoretical issues coupling with the hydrodynamics equations [24], but also in practical approaches in recent studies. The density field was determined by the indicator function derived from an irrotational discrete delta vector in the droplet simulations with a large density ratio in [25]. To incorporate the viscoelasticity effect of Oldroyd-B fluid, the smoothed Heaviside indicator function was used for simulating the dynamics of Newtonian vesicle in [26]. The role of the indicator function was represented by the solid fraction in the diffuse-interface IB framework for conjugate-heattransfer problems proposed by [27]. Moreover, the indicator function was used to identify the spatial distribution of physical quantities to predict the structure and refractive index profile of fiber-optic components in [28]. In [21], the authors presented the indicator function $I(x, y)$ based on a continuous version as follows:

$$
\begin{equation*}
I(x, y)=\int_{A} \delta\left(x-x^{\prime}\right) \delta\left(y-y^{\prime}\right) d a^{\prime} \tag{1}
\end{equation*}
$$

where $A$ is an area. Then, symbolically, we have the following equation:

$$
\begin{align*}
\nabla I(x, y) & =\int_{A} \nabla\left(x-x^{\prime}\right) \delta\left(y-y^{\prime}\right) d a^{\prime} \\
& =-\int_{A} \nabla^{\prime} \delta\left(x-x^{\prime}\right) \delta\left(y-y^{\prime}\right) d a  \tag{2}\\
& =-\oint_{\partial A} \delta\left(x-x^{\prime}\right) \delta\left(y-y^{\prime}\right) n^{\prime} d a^{\prime} \\
& =-\int_{\partial A} \delta\left(x-x^{\prime}\right) \delta\left(y-y^{\prime}\right) n^{\prime} d a^{\prime}
\end{align*}
$$

where $n^{\prime}$ is the outward unit normal vector at the domain boundary $\partial A$.

However, it is not straightforward to connect the relationship between the numerical scheme and the continuous version of the indicator function. Therefore, the objective of this paper is to present a detailed derivation and numerical investigation of the indicator function for front-tracking methods.

The paper is organized in the following manner. In Section 2, we present the detailed derivation of the indicator function. In Section 3, various numerical experiments are performed to investigate the effect of parameters such as distance between points, uniformity of the distance, and types of the Dirac delta functions on the indicator function. In Section 5, conclusions are drawn.

## 2. Derivation of the Indicator Function

Let $\delta(x)$ be a smoothed Dirac delta function and satisfy the following equation:


Figure 1: Schematic of delta function $\delta(x)$ and shifted delta function $\delta\left(x-X_{0}\right)$.

$$
\begin{equation*}
\int_{-\infty}^{\infty} \delta\left(x-x_{0}\right) d x=1 \text { for any } x_{0} \in \mathbb{R} \tag{3}
\end{equation*}
$$

For example, a 4-point delta function [29-32] is given as follows:

$$
\begin{equation*}
\delta(x)=\frac{1}{h} \phi\left(\frac{x}{h}\right), \tag{4}
\end{equation*}
$$

where $h$ is the space step size in the discrete equation and given by the following equation:

$$
\phi(r)= \begin{cases}\frac{\left(3-2|r|+\sqrt{1+4|r|-4 r^{2}}\right)}{8}, & \text { if }|r| \leq 1,  \tag{5}\\ \frac{\left(5-2|r|-\sqrt{-7+12|r|-4 r^{2}}\right)}{8}, & \text { if } 1 \leq|r| \leq 2, \\ 0, & \text { otherwise }\end{cases}
$$

Now, we consider a series of the shifted delta function $\delta\left(x-X_{0}\right)$ for some $-h<X_{0}<h$ and it is shown in Figure 1 for some $X_{0}=0.7$ and $h=1$, here we take $h=1$ for better visualization. We describe the $h$ effect in Section 3.1.

Let $x_{i}=i h$ for $i= \pm 1, \pm 2, \ldots$. From the definition (4), we have the following equation:

$$
\begin{align*}
\sum_{i=-\infty}^{\infty} h \delta\left(x_{i}-X_{0}\right)= & h \delta\left(x_{-2}-X_{0}\right)+h \delta\left(x_{-1}-X_{0}\right) \\
& +h \delta\left(x_{0}-X_{0}\right)+h \delta\left(x_{1}-X_{0}\right)  \tag{6}\\
& +h \delta\left(x_{2}-X_{0}\right)=1
\end{align*}
$$

For any $X_{0} \in \mathbb{R}, \sum_{i=-\infty}^{\infty} h \delta\left(x_{i}-X_{0}\right)=1$ holds.

Let us consider a point $x=X_{0}$ at the interface in the onedimensional space. We want to construct an indicator function which increases monotonically from zero to one across the point $X_{0}$ as $x$ changes its position from left to right of the point, i.e.,

$$
\begin{equation*}
I(x)=\int_{-\infty}^{x} \delta\left(t-X_{0}\right) d t \tag{7}
\end{equation*}
$$

Figure 2 shows the delta function $\delta(x)$ (solid line) and the indicator function $I(x)$ (dashed line) with $h=1$.

Taking a second derivative of equation (7) with respect to $x$, then we have the following equation:

$$
\begin{equation*}
I_{x x}(x)=\delta_{x}\left(x-X_{0}\right) \tag{8}
\end{equation*}
$$

It implies that the indicator function can be found by solving the Poisson equation (8). Let us consider a 4 -point delta function with $X_{0}=0$ and a discretization for equation (8) as follows:

$$
\begin{equation*}
I_{i+1}-2 I_{i}+I_{i-1}=\frac{\left(\delta_{i-1}-\delta_{i+1}\right)}{2} \tag{9}
\end{equation*}
$$

where we have used the unit space step size $h=1$ and $I_{i}=I\left(x_{i}\right)=I(i), \delta_{i}=\delta\left(x_{i}\right)=\delta(i)$. Let $I_{-4}=0$ and $I_{4}=I_{3}$, then we can rewrite equation (9) for $i=-3,-2, \ldots, 3$ as a matrix vector format:

$$
\left(\begin{array}{ccccccc}
-2 & 1 & 0 & 0 & 0 & 0 & 0  \tag{10}\\
1 & -2 & 1 & 0 & 0 & 0 & 0 \\
0 & 1 & -2 & 1 & 0 & 0 & 0 \\
0 & 0 & 1 & -2 & 1 & 0 & 0 \\
0 & 0 & 0 & 1 & -2 & 1 & 0 \\
0 & 0 & 0 & 0 & 1 & -2 & 1 \\
0 & 0 & 0 & 0 & 0 & 1 & -1
\end{array}\right)\left(\begin{array}{l}
I_{-3} \\
I_{-2} \\
I_{-1} \\
I_{0} \\
I_{1} \\
\frac{\delta_{0}-\delta_{-2}}{2} \\
I_{3} \\
\frac{\delta_{-1}-\delta_{-3}}{2} \\
\frac{\delta_{1}-\delta_{-1}}{2} \\
\frac{\delta_{2}-\delta_{0}}{2} \\
\frac{\delta_{3}-\delta_{1}}{2} \\
\frac{\delta_{4}-\delta_{2}}{2}
\end{array}\right) .
$$

By multiplying both sides by the inverse of the coefficient matrix, we have the following equation:


Figure 2: Delta function $\delta(x)$ (solid line) and indicator function $I(x)$ (dashed line) with $h=1$.

$$
\left(\begin{array}{l}
I_{-3}  \tag{11}\\
I_{-2} \\
I_{-1} \\
I_{0} \\
I_{1} \\
I_{2} \\
I_{3}
\end{array}\right)=\left(\begin{array}{ccccccc}
-1 & -1 & -1 & -1 & -1 & -1 & -1 \\
-1 & -2 & -2 & -2 & -2 & -2 & -2 \\
-1 & -2 & -3 & -3 & -3 & -3 & -3 \\
-1 & -2 & -3 & -4 & -4 & -4 & -4 \\
-1 & -2 & -3 & -4 & -5 & -5 & -5 \\
-1 & -2 & -3 & -4 & -5 & -6 & -6 \\
-1 & -2 & -3 & -4 & -5 & -6 & -7
\end{array}\right)\left(\begin{array}{l}
\frac{\delta_{-2}-\delta_{-4}}{2} \\
\frac{\delta_{-1}-\delta_{-3}}{2} \\
\frac{\delta_{0}-\delta_{-2}}{2} \\
\frac{\delta_{1}-\delta_{-1}}{2} \\
\frac{\delta_{3}-\delta_{0}}{2} \\
\frac{\delta_{4}-\delta_{2}}{2}
\end{array}\right) .
$$

From equation (11), we obtain the following results:

$$
\begin{align*}
I_{-3} & =\frac{1}{2}\left(\delta_{-4}+\delta_{-3}-\delta_{3}-\delta_{4}\right)=0 \\
I_{i} & =\frac{1}{2} \delta_{i}+\sum_{j=-2}^{i-1} \delta_{j} \text { fori } i=-2,-1,0,1,2  \tag{12}\\
I_{3} & =\delta_{-2}+\delta_{-1}+\delta_{0}+\delta_{1}+\delta_{2}=1
\end{align*}
$$

where we have used $\delta_{-4}=\delta_{-3}=\delta_{3}=\delta_{4}=0$. Since $\delta(x) \geq 0$, the indicator function $I(x)$ is a monotonically increasing


Figure 3: (a) Delta function $\delta(x)$ (solid line) and indicator function $I(x)$ (dashed line) with $X_{0}=0$. (b) Discrete delta and indicator functions with various point positions, $-1<X_{0}<1$.


Figure 4: Sum of the Dirac delta functions $\delta(x+4)-\delta(x-4)$ and the associated indicator function $I(x)$.
function, i.e., $0=I_{-4}=I_{-3} \leq I_{-2} \leq \ldots \leq I_{3}=I_{4}=1$. Figure 3(a) shows discrete delta function $\delta_{-4}=\delta_{-3}=\delta_{-2}=$ $\delta_{2}=\delta_{3}=\delta_{4}=0, \delta_{-1}=\delta_{1}=0.25, \delta_{0}=0.5$ and discrete indicator function $I_{-4}=I_{-3}=I_{-2}=0, I_{-1}=0.125, I_{0}=0.5$, $I_{1}=0.975, I_{2}=I_{3}=I_{4}=1$ obtained from equation (12). Figure 3(b) shows discrete delta and indicator functions with various point positions, $-1<X_{0}<1$.

Next, let us consider a hat shaped indicator function as shown in Figure 4 and the sum of the corresponding Dirac delta functions is $\delta(x+4)-\delta(x-4)$. If we solve the following equation for the indication function with zero Dirichlet boundary condition, then we have the result in Figure 4.

$$
\begin{equation*}
I_{x x}(x)=\delta_{x}(x+4)-\delta_{x}(x-4) \tag{13}
\end{equation*}
$$



Figure 5: Eulerian points $\mathbf{x}_{i j}$, Lagrangian points $\mathbf{X}_{l}$, and normal vector $\mathbf{n}$ for Lagrangian points in the computational domain $\Omega=(a, b) \times(c, d)$.

Now, we consider this in the two-dimensional space. Let a computational domain $\Omega=(a, b) \times(c, d)$ be partitioned in Cartesian geometry. Let $N_{x}$ and $N_{y}$ be the number of cells in the $x$ - and $y$-directions, respectively. We assume a uniform mesh with mesh spacing $h=(b-a) / N_{x}=$ $(d-c) / N_{y}$. Let $\quad \mathbf{x}_{i j}=\left(x_{i}, y_{j}\right)=(a+i h, c+j h)$ for $i=0, \ldots, N_{x}$ and $j=0, \ldots, N_{y}$. Let $\mathbf{X}_{l}=\left(X_{l}, Y_{l}\right)$ for $l=$ $1, \ldots, M$ be a set of $M$ Lagrangian points, see Figure 5.

Because $I(\mathbf{x})$ is an indicator function, the following equation holds at the transition layer of the function:

$$
\begin{equation*}
\frac{\nabla I(\mathbf{x})}{|\nabla I(\mathbf{x})|}=-\mathbf{n}(\mathbf{x}) \tag{14}
\end{equation*}
$$

where $\mathbf{n}(\mathbf{x})$ is the outward normal vector at $\mathbf{x}$. Next, let us rewrite equation (14) as follows:

$$
\begin{equation*}
\nabla I(\mathbf{x})=-\mathbf{n}(\mathbf{x})|\nabla I(\mathbf{x})| . \tag{15}
\end{equation*}
$$

We replace $\mathbf{n}(\mathbf{x})|\nabla I(\mathbf{x})|$ in equation(14). by $\sum_{l=1}^{M} \mathbf{n}\left(\mathbf{X}_{l}\right) \delta$ $\left(\mathbf{x}-\mathbf{X}_{l}\right) \Delta s_{l}$, then we have the following equation:

$$
\begin{equation*}
\nabla I(\mathbf{x})=-\sum_{l=1}^{M} \mathbf{n}\left(\mathbf{X}_{l}\right) \delta\left(\mathbf{x}-\mathbf{X}_{l}\right) \Delta s_{l} \tag{16}
\end{equation*}
$$

where $\quad \delta\left(\mathbf{x}-\mathbf{X}_{l}\right)=\delta\left(x-X_{l}\right) \delta\left(y-Y_{l}\right) \quad$ and $\quad \Delta s_{l}=0.5$ $\left(\left|\mathbf{X}_{l+1}-\mathbf{X}_{l}\right|+\left|\mathbf{X}_{l-1}-\mathbf{X}_{l}\right|\right)$. After taking the divergence operator to equation (16), we have the following Poisson's equation:

$$
\begin{equation*}
\Delta I(\mathbf{x})=\nabla \cdot \nabla I(\mathbf{x})=-\nabla \cdot\left(\sum_{l=1}^{M} \mathbf{n}\left(\mathbf{X}_{l}\right) \delta\left(\mathbf{x}-\mathbf{X}_{l}\right) \Delta s_{l}\right) \tag{17}
\end{equation*}
$$

We solve equation (17). with appropriate Dirichlet boundary condition. In this study, we use the Gauss-Seidel method.

## 3. Numerical Experiments

Let us consider two-dimensional indicators on the computational domain, $\Omega=(-4,4) \times(-4,4)$ with $h=1$. In Figures 6 and 7, we consider uniform distance change of
$y$-direction, $\Delta s$ for fixed $X_{0}=0$. In the immersed boundary method, there are many types of delta functions. Most commonly used delta functions are 2-point [31-33], 3-point [31, 32, 34], 4-point [29-32], 4-point cosine [30, 32], and 6point $[29,31,35,36]$ functions. These functions are defined as follows:
(i) 2-point delta function [31-33]:

$$
\phi(r)= \begin{cases}1-|r|, & \text { if }|r| \leq 1  \tag{18}\\ 0, & \text { otherwise }\end{cases}
$$

(ii) 3-point delta function [31, 32, 34]:

$$
\phi(r)= \begin{cases}\frac{1}{3}\left(1+\sqrt{-3 r^{2}+1}\right), & \text { if }|r| \leq 0.5  \tag{19}\\ \frac{1}{6}\left(5-3|r|-\sqrt{-3(1-|r|)^{2}+1}\right), & \text { if } 0.5<|r| \leq 1.5 \\ 0, & \text { otherwise. }\end{cases}
$$

(iii) 4-point cosine delta function [30, 32]:

$$
\phi(r)= \begin{cases}\frac{1}{4}\left(1+\cos \left(\frac{\pi r}{2}\right)\right) & \text { if }|r| \leq 2  \tag{20}\\ 0 & \text { otherwise }\end{cases}
$$

(iv) 6 -point delta function $[29,31,35,36]$ :

$$
\phi(r)= \begin{cases}\frac{61}{112}-\frac{11}{42}|r|-\frac{11}{56}|r|^{2}+\frac{1}{12}|r|^{3} &  \tag{21}\\ +\frac{\sqrt{3}}{336}\left(243+1584|r|-748|r|^{2}-1560|r|^{3}+500|r|^{4}+336|r|^{5}-112|r|^{6}\right) & \text { if }|r| \leq 1 \\ \frac{21}{16}+\frac{7}{12}|r|-\frac{7}{8}|r|^{2}+\frac{1}{6}|r|^{3}-\frac{3}{2} \phi(|r|-1)^{1 / 2} & \text { if } 1<|r| \leq 2 \\ \frac{9}{8}-\frac{23}{12}|r|+\frac{3}{4}|r|^{2}-\frac{1}{12}|r|^{3}+\frac{1}{2} \phi(|r|) & \text { if } 2<|r| \leq 3 \\ 0 & \text { otherwise }\end{cases}
$$

Figure 6 shows 2-point, 3-point, 4-point, 4-point cosine, and 6 -point delta functions from top to bottom row with $\Delta s=1$. Figures 6(a)-6(d) are 2D delta function $\delta\left(\mathbf{x}-\mathbf{X}_{0}\right)=$ $\delta\left(x-X_{0}\right) \delta\left(y-Y_{0}\right)$ with $X_{0}=0$ and $Y_{0}=0 ; 2 \mathrm{D}$ delta functions $\delta\left(\mathbf{x}-\mathbf{X}_{l}\right)=\delta\left(x-X_{l}\right) \delta\left(y-Y_{l}\right)$ with $X_{l}=0$ and various point positions with one distance; sum of 2D delta functions $\sum_{l} \delta\left(\mathbf{x}-\mathbf{X}_{l}\right)$; and 2 D indicator function $I(\mathbf{x})$,
respectively. We obtain very similar results of indicator functions $I(\mathbf{x})$ for all cases except the 2-point delta function. Figure 7 shows the result of delta functions, the sum of delta functions, and indicator functions according to different $\Delta s=h, 2 h, 3 h$, respectively. As shown in Figure 7, in cases of $\Delta s=h$ and $\Delta s=2 h$, the indicator functions are similar. On the other hand, when $\Delta s=3 h$, the indicator function seems


Figure 6: From top to the bottom row, 2-point, 3-point, 4-point, 4-point cosine, and 6-point delta functions: (a) 2D delta function $\delta\left(\mathbf{x}-\mathbf{X}_{0}\right)=\delta\left(x-X_{0}\right) \delta\left(y-Y_{0}\right)$ with $X_{0}=0$ and $Y_{0}=0$, (b) 2 D delta functions $\delta\left(\mathbf{x}-\mathbf{X}_{l}\right)=\delta\left(x-X_{l}\right) \delta\left(y-Y_{l}\right)$ with $X_{l}=0$ and various point positions with one distance, (c) sum of 2 D delta functions $\sum_{l} \delta\left(\mathbf{x}-\mathbf{X}_{l}\right)$, and (d) 2 D indicator function $I(\mathbf{x})$.
unstable compared to Figures 7(a) and 7(b). For this reason, it is recommended to use $\Delta s$ which is at least smaller than $2 h$ for a stable indicator function construction.

Let us consider an indicator on the computational domain, $\Omega=(-15,15) \times(-15,15)$ with $h=1$ using the 4 point delta function. Figure 8 shows immersed boundary positions as a circle with different uniform distance $\Delta s, 2 \mathrm{D}$ delta functions $\delta\left(\mathbf{x}-\mathbf{X}_{l}\right)=\delta\left(x-X_{l}\right) \delta\left(y-Y_{l}\right), \quad l=1,2$, $\ldots, M$, the sum of 2 D delta functions $\sum_{l} \delta\left(\mathbf{x}-\mathbf{X}_{l}\right)$, and 2 D indicator function $I(\mathbf{x})$ from top to bottom row. We can observe that if $\Delta s$ is too large, then the result of the indicator function is not good.

The first and second columns in Figure 9 show the construction of the indicator functions with randomly
distributed points ranging $0.5 h<\Delta s_{l}<1.5 h$ and $2.5 h<\Delta s_{l}<4 h$, respectively. We can observe that the indicator function is well reconstructed from randomly distributed points ranging $0.5 h<\Delta s_{l}<1.5 h$.
3.1. Effects of $h$. In this section, we describe the effects of spatial grid size $h$. Figure 10 shows 4 -point delta functions and corresponding indicator functions for different spatial step $h=1,0.5$, and 0.25 . We observe that as $h$ is smaller, the delta function becomes sharper and the interface transition layer of the indicator function becomes shorter.

In computational domain $\Omega=(-15,15) \times(-15,15)$, we consider uniform immersed boundary points for a circle


Figure 7: From top to bottom row, 2D delta functions $\delta\left(\mathbf{x}-\mathbf{X}_{l}\right)=\delta\left(x-X_{l}\right) \delta\left(y-Y_{l}\right)$ with $X_{l}=0$ and $Y_{l}= \pm \Delta s, \pm 2 \Delta s, \ldots$, sum of 2 D delta functions $\sum_{l} \delta\left(\mathbf{x}-\mathbf{X}_{l}\right)$, and 2D indicator function $I(\mathbf{x}):(\mathrm{a}) \Delta s=h$, (b) $\Delta s=2 h$, and (c) $\Delta s=3 h$.
radius of 10 with different $h=1,0.5$, and 0.25 as shown in Figure 11(a). In Figure 11(b), we solve the indicator function, accordingly the points in Figure 11(a). It has a more stiff interface as $h$ is smaller. Figure 11(c) shows $\left\|I^{(k)}-I^{(k-1)}\right\|_{2}$ for the Gauss-Seidel iteration with a tolerance tol $=1.0 e-6$, where $k$ is the number of iterations and $I^{(0)}$ is a zero vector. We can observe that the number of iterations is getting larger as $h$ approaches zero.

## 4. Fluid Application

In this section, we give two examples for the applications of indicator function in multiphase flows and fluid-structure interaction (FSI) problem. First, we consider the two-phase flows problem as the dynamics of a droplet suspended in the simple shear flow. The droplet is between two parallel plates as shown in Figure 12. Let $\mathbf{u}(\mathbf{x}, t)=(u(\mathbf{x}, t), v(\mathbf{x}, t))$ be fluid velocity at $\mathbf{x}=(x, y)$ and time $t$. The Lagrangian variable, denoted as $\mathbf{X}(s, t)$, where $s$ is arc length parameter, expresses the immersed boundary. The governing equations on the domain $\Omega$ are as follows:

$$
\begin{align*}
\frac{\partial \mathbf{u}(\mathbf{x}, t)}{\partial t}+\mathbf{u}(\mathbf{x}, t) \cdot \nabla \mathbf{u}(\mathbf{x}, t) & =-\nabla p(\mathbf{x}, t) \\
& +\frac{1}{R e} \Delta \mathbf{u}(\mathbf{x}, t)+\frac{1}{W e} \mathbf{f}(\mathbf{x}, t)  \tag{22}\\
\nabla \cdot \mathbf{u}(\mathbf{x}, t) & =0 \\
\mathbf{f}(\mathbf{x}, t) & =\int_{\mathrm{T}} \mathbf{F}(\mathbf{X}(s, t)) \delta^{2}(\mathbf{x}-\mathbf{X}(t)) d s \\
\mathbf{F}(s, t) & =\frac{\partial^{2} \mathbf{X}(s, t)}{\partial s^{2}}, \\
\frac{\partial \mathbf{X}(t)}{\partial t} & =\mathbf{U}(\mathbf{x}(t)) \\
\mathbf{U}(\mathbf{X}(s, t)) & =\int_{\Omega} \mathbf{u}(\mathbf{x}, t) \delta^{2}(\mathbf{x}-\mathbf{X}(s, t)) d \mathbf{x} \tag{23}
\end{align*}
$$

where $\Gamma$ is the interface between the two fluids. We define the Reynolds number in equation (22) as $\operatorname{Re}=\rho_{2} R^{2} \dot{\gamma}$,


Figure 8: (a) Immersed boundary positions with different distance $\Delta s$, (b) 2D delta functions of at the boundary positions, (c) sum of 2D delta functions, and (d) 2D indicator function $I(\mathbf{x})$ with different uniform distance $\Delta s$.


FIGURE 9: (a) Randomly distributed boundary positions, (b) 2D delta functions of boundary positions, (c) sum of 2D delta functions, and (d) 2D indicator function $I(\mathbf{x})$ with different nonuniform distance $\Delta s_{l}$.


Figure 10: Delta and indicator functions with different spatial steps $h=1,0.5$, and 0.25 .
$W e=\operatorname{Re} \cdot C a$, and $C a=\eta R \dot{\gamma} / \sigma$, where $\rho$ is fluid density, $\eta$ fluid viscosity, $R$ droplet radius, $\dot{\gamma}$ the shear rate, and $\sigma$ the surface tension coefficient. We define a gradient field,

$$
\begin{equation*}
\nabla I(\mathbf{x}, t)=\int_{\Gamma} \mathbf{n}(\mathbf{X}(s, t)) \delta^{2}(\mathbf{x}-\mathbf{X}(s, t)) d s \tag{24}
\end{equation*}
$$

To compute the indicator function, the following equation is solved:

$$
\begin{equation*}
\Delta I(\mathbf{x}, t)=\nabla \cdot \int_{\Gamma} \mathbf{n}(\mathbf{X}(s, t)) \delta^{2}(\mathbf{x}-\mathbf{X}(s, t)) d s \tag{25}
\end{equation*}
$$

Then, the variable density $\rho$ and viscosity $\mu$ are defined as follows:

$$
\begin{align*}
& \rho(I(\mathbf{x}, t))=\rho_{2}+\left(\rho_{1}-\rho_{2}\right) I(\mathbf{x}, t) \\
& \mu(I(\mathbf{x}, t))=\mu_{2}+\left(\mu_{1}-\mu_{2}\right) I(\mathbf{x}, t) \tag{26}
\end{align*}
$$

where $\rho_{i}$ and $\mu_{i}$ for $i=1,2$ are density and viscosity of fluid $i$, respectively.

For simplicity, the viscosity and density of the droplet is set equal to that of the matrix phase. Our focus is on the generations of indicator functions from the unit normal of the interface in the drop deformation process. We refer to [37] for the details of numerical algorithms. We consider the simulations of drop deformation on $\Omega=[0,2] \times[0,2]$ with a space mesh size $128 \times 128$. A droplet with a radius 0.5 is put into the center of the domain with initially three sets of different Lagrangian points for comparisons. Delta function (5), Reynolds number $\mathrm{Re}=1$, and capillary number $\mathrm{Ca}=$ 0.3 are used in simulations. We compared three different $\Delta s$ settings to investigate the effect of $\Delta s$ on the formation of indicator function at the end time $t=2.5$. The initial $\Delta s_{i n i}$ values are set to $1.0 \mathrm{~h}, 5.0 \mathrm{~h}$ and 10.0 h with the number of boundary points 201,38 , and 20 , separately, see the first column of Figure 13. While the drop deforms, the interface is remeshed by adding and deleting points with the condition $\Delta s_{\max } / \Delta s_{\min } \leq 2.0$ for different $\Delta s_{\text {ini }}$ leading to the boundary points being more evenly distributed. As the time evolves,
the indicator function could be well resolved for $\Delta s_{\text {ini }}=1.0 \mathrm{~h}$ as shown in the second row of Figure 13. However, increasing $\Delta s$ will lead the indicator function a slightly serrated border as shown in the second and sixth rows of Figure 13 (case $\Delta s_{\text {ini }}=5.0 \mathrm{~h}$ and 10.0 h ), which is unacceptable to separate two immiscible fluids. Moreover, we could see the drop interface shrink in the normal direction quickly by setting the $\Delta s$ larger while drop deforms in the shear flow, as shown in the odd row of Figure 13. Table 1 lists the area losses as the drop evolves at the time $t=2.5$. It shows that the larger $\Delta s$, the more the drop area loss.

Next, we take the direct-forcing immersed boundary method as an example application in the FSI problem of the indicator function. The nondimensional governing equations are similar to the multiphase flows and read as follows:

$$
\begin{align*}
& \frac{\partial \mathbf{u}(\mathbf{x}, t)}{\partial t}+\mathbf{u}(\mathbf{x}, t) \cdot \nabla \mathbf{u}(\mathbf{x}, t)=-\nabla p(\mathbf{x}, t) \\
&+\frac{1}{R e} \Delta \mathbf{u}(\mathbf{x}, t)+\mathbf{f}(\mathbf{x}, t), \\
& \nabla \cdot \mathbf{u}(\mathbf{x}, t)= 0, \\
& \mathbf{f}(\mathbf{x}, t)= \int_{\mathrm{T}} \mathbf{F}(\mathbf{X}(s)) \delta^{2}(\mathbf{x}-\mathbf{X}(t)) d s, \\
& \frac{\partial \mathbf{X}(s)}{\partial t}= \mathbf{U}(\mathbf{X}(s)), \\
& \mathbf{U}(\mathbf{X}(s))=\int_{\Omega} \mathbf{u}(\mathbf{x}, t) \delta^{2}(\mathbf{x}-\mathbf{X}(t)) d \mathbf{x}, \tag{27}
\end{align*}
$$

where $\mathbf{u}$ is the velocity, $p$ is the pressure, $\mathbf{u}$ is the body force, and Re is the Reynolds number. From equation (27), we can get the following equation:


(a)


(b)


$$
\begin{aligned}
& \text { - } h=1 \\
& \text { * } h=0.5 \\
& \square-h=0.25
\end{aligned}
$$

(c)

Figure 11: (a)-(b) From left to right column, different spatial steps $h=0.25,0.5$, and 1 : (a) Lagrangian point $\mathbf{X}_{l}$ and (b) indicator functions. (c) $l_{2}$-error of $I^{(k)}$ and $I^{(k-1)}$ for the number of iterations $k$ using the gauss-seidel method with different spatial steps $h=1,0.5$, and 0.25 .

$$
\begin{gather*}
\mathbf{f}=\frac{\partial \mathbf{u}}{\partial t}+\mathbf{u} \cdot \nabla \mathbf{u}+\nabla p-\frac{1}{R e} \Delta \mathbf{u}=\frac{\partial \mathbf{u}}{\partial t}-\mathbf{r h s}, \\
\frac{\partial \mathbf{u}}{\partial t}-\mathbf{r h s}=\frac{\mathbf{u}^{n+1}-\mathbf{u}^{n}}{\Delta t}-\mathbf{r h s} \tag{28}
\end{gather*}
$$

where $n$ and $n+1$ are time steps, and rhs = $-(\mathbf{u} \cdot \nabla \mathbf{u}+\nabla p-1 / \operatorname{Re} \Delta \mathbf{u})$. The forcing exerted on the

Lagrangian points at the immersed boundary can be written as follows:

$$
\begin{equation*}
\mathbf{F}(\mathbf{X}(s))=\frac{\mathbf{U}^{n+1}(\mathbf{X}(t))-\mathbf{U}^{*}(\mathbf{X}(s))}{\Delta t} \tag{29}
\end{equation*}
$$

where $\mathbf{U}^{n+1}(\mathbf{X}(s))$ is the desired velocity and $\mathbf{U}^{*}(\mathbf{X}(s))$ is a temporary velocity of the boundary. The algorithm in detail can be found in [38]. The difference of immersed boundary


Figure 12: Schematic illustration of a liquid droplet.
to handle the multiphase flows and FSI lies on the how to evaluate the boundary force.

For a moving rigid body immersed on the fluid, the imposed force term, the authors in [39] pointed out it
consists of two parts that can be written as the following equations (30) and (31):

$$
\begin{align*}
& f_{x}=\oint_{\partial \Omega} F_{x}(X(s)) d s-\iint_{\Omega} \frac{\partial u}{\partial t} d V \approx \sum_{j=1}^{m} F_{x}\left(X_{j}\right) \Delta V_{j}-\sum_{i=1}^{N_{x}} \sum_{j=1}^{N_{y}} \frac{u_{i j}^{n+1}-u_{i j}^{n}}{\Delta t} I_{i j} \Delta V_{i j}  \tag{30}\\
& f_{y}=\oint_{\partial \Omega} F_{y}(X(s)) d s-\iint_{\Omega} \frac{\partial v}{\partial t} d V \approx \sum_{j=1}^{m} F_{y}\left(X_{j}\right) \Delta V_{j}-\sum_{i=1}^{N_{x}} \sum_{j=1}^{N_{y}} \frac{v_{i j}^{n+1}-v_{i j}^{n}}{\Delta t} I_{i j} \Delta V_{i j} \tag{31}
\end{align*}
$$

where $\partial \Omega$ is the solid boundary, $\Omega$ the solid phase, $\Delta t$ the time step, $I_{i j}$ the indicator function, $u_{i j}, v_{i j}$ the velocity in Eulerian gird, $N_{x}, N_{y}$ the mesh grid numbers, and $V_{i j}$ the mesh volume in $i j$ grid point.

One force in the first terms of right hand sides of equations (30) and (31) is that of a submersed body acting on the external fluid, which can be determined simply by integrating the immersed boundary force densities by equation (29) on the solid boundary that is already evaluated by the direct forcing. Another one that contributes to the unsteady flow inside the solid phase, is the so-called internal or virtual flow force [39] in the second terms of the right hand side of equations (30) and (31). Usually, if the object is fixed in the fluid, we impose the $\mathbf{f}=\left(f_{x}, f_{y}\right)$ term to satisfy the no-slip boundary condition on the solid surface leading to the second term equals zero. However, when the solid boundary is moving, we should not neglect the effect of the virtual flow, even though its contribution is small [39]. For a complexly enclosed geometry, the indicator function can be used to calculate to its inside virtual flow by letting its value 1 , and outside 0 of the moving solid object.

For this moving solid problem in fluid, we consider a 2D NACA0012 airfoil as the representative of a fish body
periodically undulating in the incoming flow. The study of this kind of fish kinematics can be found in [40, 41]. Our focus is on the fish shape expressed by the indicator function. As in the first example of drop deformation in the shear flow, we investigate the effect of $d s$ on the indicator function. The computational domain is $[0,4] \times[0,2]$, with mesh size 256 for $y$ direction. The airfoil is initially represented by 202,102 , and 42 Lagrangian points leading to the point distance $\Delta s$ is about $0.918 h, 2.117 h$, and $6.134 h$, where $h$ is the Eulerian grid size. The odd rows of Figure 14 shows the fish undulations with outer normal vectors attached on the boundary points for different $\Delta s$ settings. Because the boundary deforms not too much, we do not need to remesh it. As shown in Figure 14, by increasing $\Delta s$, the indicator functions get zigzag distribution on their boundaries between 0 and 1 , which could lead to the calculations of $f_{x}$ and $f_{y}$ in equations (30) and (31) a little oscillations. However, the $\Delta s \approx 1.0 h$ is good enough to resolve the smooth distribution of indicator function as in the first drop deformation example. It should be noted that the diffused fish interface is not well resolved at the high curvature section of the fishtail for larger $\Delta s$ as shown in the third case of Figure 14. Letting $\Delta s \approx 1.0 h$ resolves it quite well.


Figure 13: From top to bottom, the first, third, and fifth rows are the boundary interfaces with different initial $\Delta s_{\text {ini }}=1.0 h, 5.0 h$ and 10.0 h (attached with normal vectors). The even rows are corresponding to theirs' indicator functions. Column (a)-(c), drop evolutions at different times.


Figure 14: From top to bottom, the first, third, and fifth rows are fish boundaries with different initial $\Delta s=0.918 h, 2.117 h$ and $6.134 h$ (attached with normal vectors), corresponding to 202, 102, and 42 boundary points. The even rows are corresponding to theirs' indicator functions. Column (a)-(d), fish undulates at different times in a period. (a) $t=0$. (b) $t=0.15$. (c) $t=0.3$. (d) $t=0.45$.

Table 1: Drop area loss for the different $\Delta s_{\text {ini }}$ at time $t=2.5$.

| $\Delta s_{\text {ini }}$ | 1.0 h | 5.0 h | 10.0 h |
| :--- | :---: | :---: | :---: |
| Initial area | 0.7852 | 0.7822 | 0.7725 |
| Final area | 0.7852 | 0.6856 | 0.2070 |
| Area loss (\%) | 0.0 | 12.35 | 73.20 |

## 5. Conclusion

In this article, we presented a detailed derivation and numerical investigation of an indicator function together with its fluid applications in the immersed boundary method. We used the discrete Dirac delta function to construct an indicator function from a set of Lagrangian points. We solved the resulting discrete Poisson equation with the zero Dirichlet boundary condition using an iterative method. We
presented several computational tests to study the effect of parameters such as distance between points, uniformity of the distance, and types of the Dirac delta functions on the indicator function. Two applications indicate that the construction of the indicator function are effective, multipurpose, and easy-to-use in separating the changing phases. We note that the indicator function can be extended directly and readily to the three-dimensional space. Based on the above-given numerical investigations, as future works, we
believe that investigation of the 3D indicator function has more practical values.

## Data Availability

No data were used to support this study.

## Conflicts of Interest

The authors declare that there are no conflicts of interest.

## Acknowledgments

The study was supported by Brain Korea 21 FOUR from the Ministry of Education of the Republic of Korea.

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